

The Abel Symposium

A variational regularity theory for optimal transportation and its application to matching

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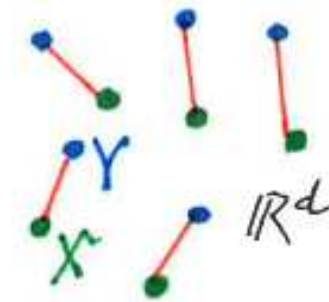
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joint work with Michael Goldman (Ann. ENS '20),
with MG & Martin Huesmann (CPAM '23),
with MH & Francesco Mattesini (submitted PTRF).

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Optimal matching of random point clouds

Matching of two locally finite point clouds $\{X\}$ and $\{Y\} \subset \mathbb{R}^d$, amounts to a pairing $\{(X, Y)\} \subset \mathbb{R}^d \times \mathbb{R}^d$.



Optimality means **cyclical monotonicity**

\forall matched finite subset $\{X_1, \dots, X_N\}$
 $\{Y_1, \dots, Y_N = Y_0\}$.

$$\sum_{n=1}^N |Y_n - X_n|^2$$

$$\leq \sum_{n=1}^N |Y_{n-1} - X_n|^2$$

$$\iff \sum_{n=1}^N (Y_n - Y_{n-1}) \cdot X_n \geq 0.$$

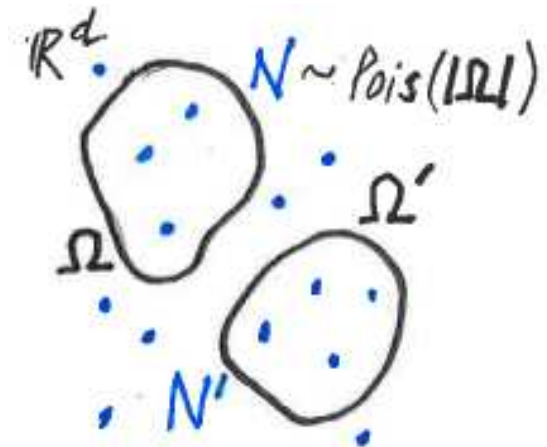
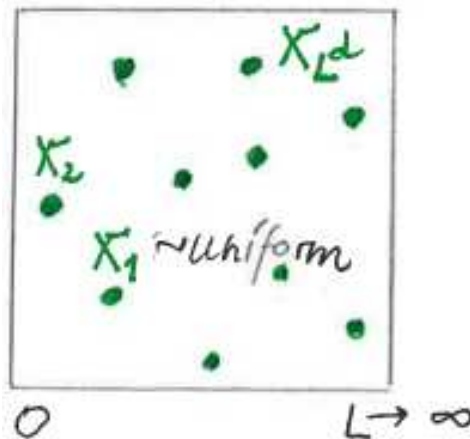
Simplest *random* setting:

$\{X\}, \{Y\}$ indep. samples of **Poisson point processes**.

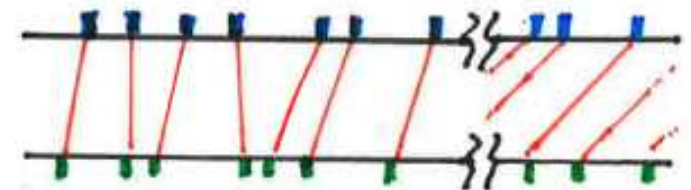
The Poisson point process

Locally finite point cloud via Poisson point process of unit intensity (means that distance between points $O(1)$)

canonical vs.
grand-canonical
definition



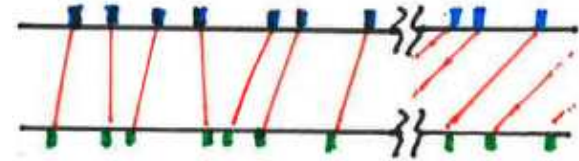
Seek cyclically monotone matching of two independent Poisson point processes
– mean distance diverges in $d = 1$.



Matching depends on dimension d ...

Cyclically monotone matching of two independent Poisson point processes

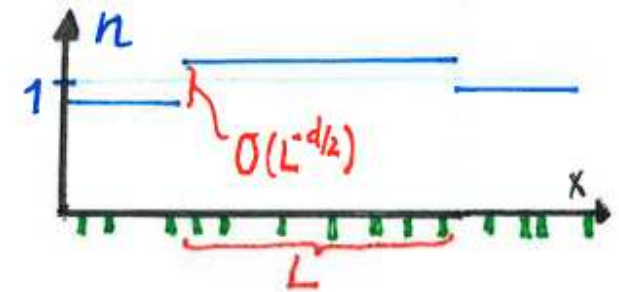
– distances diverge like square-root for $d = 1$.



Cause of divergence:
mesoscopic fluctuations

of number density n

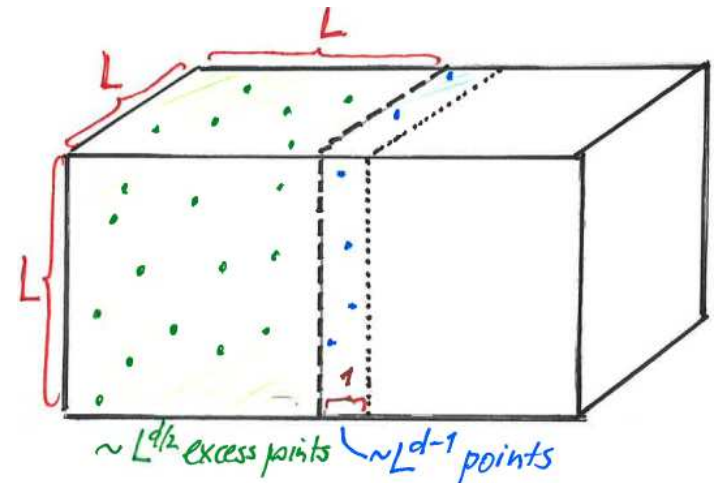
$= O(\frac{1}{\sqrt{L^d}})$; lower for higher d .



Number of excess points $= O(\sqrt{L^d})$,

number of points in

(width one) boundary layer $= O(L^{d-1})$.



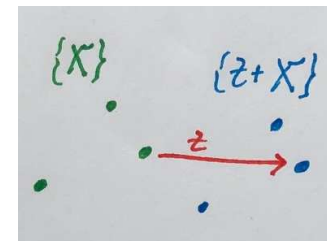
... critical dimension $d = 2$

Impose statistical translation invariance (“Stationarity”) of matching

Cyclically monotone matching of two independent Poisson point processes $\{X\}$, $\{Y\}$.

Poisson point process is *stationary*:

\forall shift vectors $z \in \mathbb{R}^d$ $\{z + X\} =_{\text{law}} \{X\}$.



Seek random point cloud $\{(X, Y)\}$ in $\mathbb{R}^d \times \mathbb{R}^d$ s. t. marginals are independent Poisson point processes, coupling is cyclically monotone almost surely, and $\forall z \in \mathbb{R}^d$ $\{(z + X, z + Y)\} =_{\text{law}} \{(X, Y)\}$.

Critical dimension $d = 2$ rigorously captured

Interest in Combinatorics (eg. Ajtai et al. '84),
Probability Theory (Talagrand '92+, Holroyd-Peres '11+),
Physics (eg. Parisi et al. '14),
Analysis (eg. Ambrosio et al. '16+).

Seek random point cloud $\{(X, Y)\}$ in $\mathbb{R}^d \times \mathbb{R}^d$ s. t.
marginals are independent Poisson point processes,
coupling is cyclically monotone almost surely,
and $\forall z \in \mathbb{R}^d \quad \{(z + X, z + Y)\} =_{\text{law}} \{(X, Y)\}$.

Theorem (Huesmann&Sturm '13)

For $d > 2$ have existence.

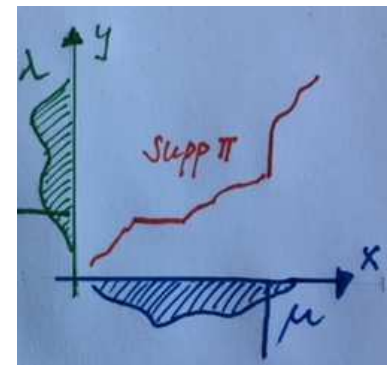
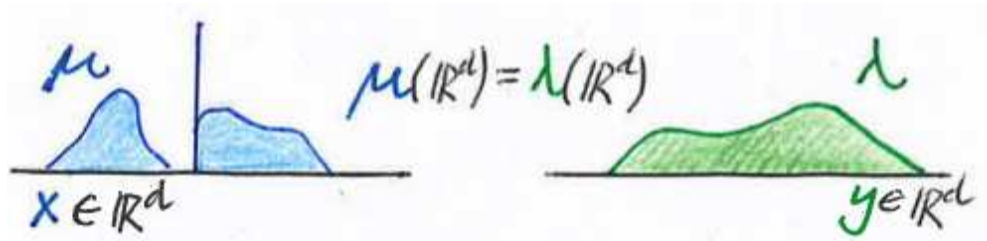
Theorem (H.&Mattesini&0. '21)

For $d \leq 2$ have non-existence.

Proof via continuum/PDE perspective

Optimal Transportation in Kantorowicz' formulation

Given two measures



seek transfer plan π , i. e. $\pi(U \times \mathbb{R}^d) = \mu(U)$, $\pi(\mathbb{R}^d \times V) = \lambda(V)$
that minimizes Euclidean transport cost $\int_{\mathbb{R}^d \times \mathbb{R}^d} |y-x|^2 \pi(dx dy)$.

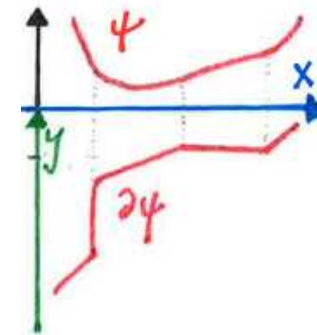
Minimum $=: W_2^2(\mu, \lambda)$ (squared) Wasserstein distance.

From optimal transportation to Monge-Ampère

Minimize $\int_{\mathbb{R}^d \times \mathbb{R}^d} |y - x|^2 \pi(dx dy)$
among all $\pi(dx dy)$ with marginals $\mu(dx)$ and dy .

Support of optimal transfer plan π
is cyclically monotone;
hence \exists convex ψ

$\text{supp } \pi \subset \{ (x, y) \mid y \in \text{sub-gradient } \partial\psi(x) \}$.



\forall test functions ζ $\int \zeta(\nabla\psi(x)) \mu(dx) = \int \zeta(y) dy$.

In smooth case, this amounts to $\det D^2\psi = \mu$,
an instance of the Monge-Ampère equation.

Nature of the Monge-Ampère equation

Recall Monge-Ampère: $\det D^2\psi = 1$.

Fully non-linear with $F(A) := \det A - 1$.

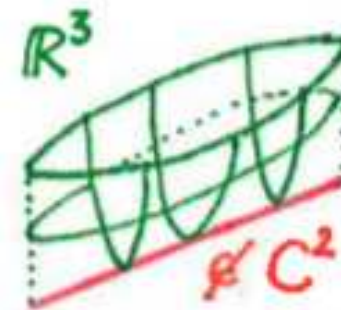
However elliptic: $F(A) > F(A')$ for $A > A' \geq 0$;
satisfies comparison principle.

However degenerate: \leftrightarrow affine invariant (non-compact $SL(d)$).

Cf. Laplacian $F(A) = \text{tr} A$: rotation invariant (compact $SO(d)$).

Caffarelli's '90 breakthrough:

comparison principle,
affine invariance,
compactness.



Pogorelov's
example is
worst case

Monge-Ampère equation at crossroads
of fully nonlinear and variational.

Parisi's heuristics for semi-discrete matching

$\mu = \sum_{x \in X} \delta_x$ Poisson, $\lambda = dy$ Lebesgue

$\det D^2 \psi - 1 = \mu - 1 \approx_{\text{law}} \xi$ on scales $R \gg 1$, $\xi = \text{white noise}$,
meaning $\int \eta_R(\mu - 1) \approx_{\text{law}} \int \eta_R \xi$, with $\eta_R(x) = \frac{1}{R^d} \eta(\frac{x}{R})$, $\eta \in C_0^\infty$ fixed.

Consider $\Delta \phi = \xi$ (so that $\nabla \phi$ behaves as Gaussian free field)
 $\xi \sim R^{-\frac{d}{2}} \ll 1$ on scales $R \gg 1 \implies D^2 \phi \ll 1$ on scales $\gg 1 \implies$
 $\det(D^2(\phi + \frac{1}{2}|x|^2)) - 1 = \det(D^2 \phi + \text{id}) - 1 \approx \text{tr} D^2 \phi = \xi$
on scales $\gg 1$.

Hence expect $\nabla \psi - x \approx_{\text{law}} \nabla \phi$ on scales $\gg 1$;

in terms of "displacement" $(y-x)\pi(dx dy)$:

$$\frac{\int \eta_R(x) (y-x) \pi(dx dy)}{\int \eta_R(x) \pi(dx dy)} \approx_{\text{law}} \int \eta_R \nabla \phi.$$

Quantitative large-scale linearization

Compare: $\pi(dx dy)$ optimal for $\mu(dx)$ and dy
with $-\Delta\phi = \mu - 1$.

Rate function $D(R) \uparrow$, $\frac{D(R)}{R} \downarrow$
in strengthened sense of Dini $\sum_{k=0}^{\infty} \frac{D(2^k R)}{2^k R} \leq 4 \frac{D(R)}{R}$.

Theorem 1 (Goldman&Huesmann&O.) Provided

$$\frac{1}{|B_R|} W_{B_R}^2(\mu, \kappa_R) + R^2(\kappa_R - 1)^2 \leq D(R) \quad \text{for all } R \geq 1,$$

$$\text{then } \left| \frac{\int \eta_R(x)(y-x)\pi(dx dy)}{\int \eta_R(x)\pi(dx dy)} - \int \eta_R \nabla \phi \right| \lesssim \frac{D(R)}{R} \quad \text{for all } R \geq 1.$$

Confirms (deterministic part of) Parisi's heuristics.

Relies on large-scale regularity theory.

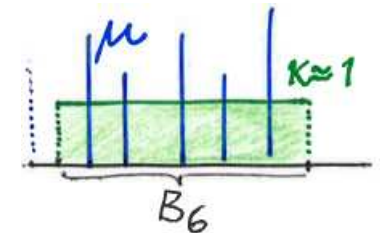
Large-scale regularity via harmonic approximation

Consider $\pi(dx dy)$ optimal for $\mu(dx)$ and $\lambda(dy)$.

Local energy
$$E := \int_{(B_6 \times \mathbb{R}^d) \cup (\mathbb{R}^d \times B_6)} |y - x|^2 \pi(dx dy),$$

Local data size²
$$D := W_{B_6}^2(\mu, \kappa = \text{const}) + (\kappa - 1)^2$$

 + same for λ



Theorem 2 (Goldman&Huesmann&O.)

$$\forall \tau > 0 \quad \exists \epsilon(\tau, d) > 0, C(\tau, d) < \infty \quad \text{s. t.} \quad E + D \leq \epsilon \implies$$

$$\exists \nabla \phi \text{ harmonic,} \quad \int_{B_1} |\nabla \phi|^2 \leq C(E + D),$$

$$\int_{(B_1 \times \mathbb{R}^d) \cup (\mathbb{R}^d \times B_1)} |y - x - \nabla \phi(y)|^2 \pi(dx dy) \leq \tau E + CD.$$

Harmonic approximation: correct homogeneities ...

$$E := \int_{(B_6 \times \mathbb{R}^d) \cup (\mathbb{R}^d \times B_6)} |y - x|^2 \pi(dx dy), \text{ quadratic in solution,}$$

$$D := W_{B_6}^2(\mu, \kappa = \text{const}) + (\kappa - 1)^2 + \text{same for } \lambda, \text{ quadratic in data.}$$

Theorem 2

$$\forall \tau > 0 \quad \exists \epsilon(\tau, d) > 0, C(\tau, d) < \infty \quad \text{s. t.} \quad E + D \leq \epsilon \implies$$

$$\exists \nabla \phi \text{ harmonic,} \quad \int_{B_1} |\nabla \phi|^2 \leq C(E + D),$$

$$\int_{(B_1 \times \mathbb{R}^d) \cup (\mathbb{R}^d \times B_1)} |y - x - \nabla \phi(y)|^2 \pi(dx dy) \leq \tau E + CD.$$

Compare to $\int_{B_1} L(\nabla u - \nabla \phi) \leq \tau \int_{B_6} L(\nabla u) + C \int_{B_6} |f|^2$

for $-\nabla \cdot DL(\nabla u) = \nabla \cdot f$ with uniformly convex L .

... and correct metric

Definition of $\nabla\phi$ via Neumann-Poisson problem

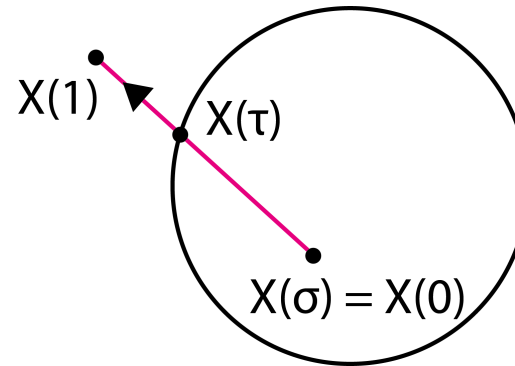
Consider trajectories $[0, 1] \ni t \mapsto X(t) := ty + (1 - t)x$.

Consider when & where trajectories $\begin{pmatrix} \text{enter} \\ \text{exit} \end{pmatrix} B_R$:

$$\begin{pmatrix} \sigma \\ \tau \end{pmatrix} := \begin{pmatrix} \min \\ \max \end{pmatrix} \{ t \in [0, 1] \mid X(t) \in \bar{B}_R \},$$

$$\int \zeta dg = \int \zeta(X(\sigma))\pi(dxdy),$$

$$\int \zeta df = \int \zeta(X(\tau))\pi(dxdy).$$



Define $\nabla\phi$ by the Neumann-Poisson problem:

$$-\Delta\phi = \frac{(\lambda - \mu)(B_R)}{|B_R|} \text{ in } B_R \quad \text{and} \quad \nu \cdot \nabla\phi = f - g \text{ on } \partial B_R.$$

Optimize in $R \in [2, 3]$.

A numerical illustration in case of matching

Matching of $\{X\}$ and $\{Y\}$; consider a square Q



Define $\nabla\phi$ by the Neumann-Poisson problem:

$$-\Delta\phi = \sum_{Y \in Q} \delta_Y - \sum_{X \in Q} \delta_X \quad \text{in } Q,$$

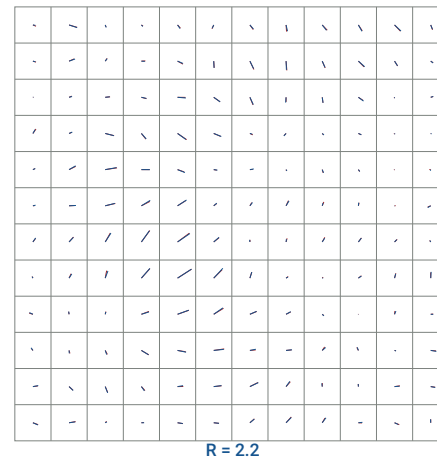
$$\nu \cdot \nabla\phi = \sum_{Y \in Q, X \notin Q} \delta_Z - \sum_{Y \notin Q, X \in Q} \delta_Z \quad \text{on } \partial Q.$$

A numerical illustration in case of matching

Recall $-\Delta\phi = \sum_{Y \in Q} \delta_Y - \sum_{X \in Q} \delta_X$ in Q ,

$\nu \cdot \nabla\phi = \sum_{Y \in Q, X \notin Q} \delta_Z - \sum_{Y \notin Q, X \in Q} \delta_Z$ on ∂Q .

Check $\nabla\phi \approx Y - X$ when averaged on scale R ;



R larger – better agreement; side-length of $Q = 75$.

Finite Element discretization, simulations by R. Kriemann

Analogies to minimal surfaces (Schoen&Simon '82)

Approximate minimal surface by harmonic graph /
approximate displacement by harmonic gradient.

Use: Object is minimizing under compact perturbations.

Don't use: Euler-Lagrange equation (= first variation).

Mismatch of type of boundary condition

for construction of harmonic competitor:

graph vs. non-graph / flux vs. displacement;

choice of good radius.

Use of strict convexity to convert energy gap
into distance (“approximate orthogonality”);
need to smooth out boundary data.

Summary and outlook

Variational regularity theory for optimal transportation, mimics ϵ -regularity theory for minimal surfaces, is more robust than maximum principle-based approach, provides large-scale regularity theory for matching.

Extend non-existence result and regularity theory from $p = 2$ to $p \in (1, \infty)$ (& L. Koch), and to entropic regularization (& R. Gvalani).

Seek quantitative coupling of shot noise $\mu - 1$ with white noise ξ in $d > 2$ (Komlós&Major&Tusnády '75).